

1D CSR with shielding in the PLACET code.

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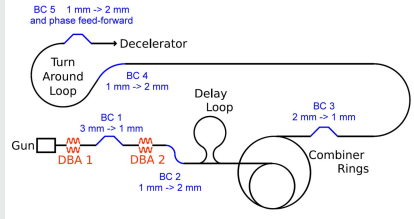
- 1 Introduction
- 2 Underlying physics
- 3 The used model
- 4 Implementing the process
- 5 Details of the interface
- 6 Examples and Benchmarking

The tracking code PLACET.



- Originally written by Daniel Schulte to simulate the beam delivery systems of future linear colliders. PLACET: Program for Linear Accelerator Correction and Efficiency Tests.
- Has since been expanded significantly with numerous contributors to simulate a multitude of new functionalities.
- E.g. 4 and 6-dim tracking, dynamic misalignments, advanced steering algorithms, fringe fields, long and short range wake fields, extension for halo generation, residual gas scattering, isr in all magnetic elements etc..
- User can set up entire beam lines with common beam line elements and track particles, extract optics functions etc.
- Interfaces Tcl, Octave and Python.
- The code uses the ultrarelativistic approximation.
- One such addition has been CSR in bending magnets (E. Adli). The implementation was in perfect agreement with Elegant.

- Emitted csr power $P \approx \frac{3^{2/3} N^2 e^2 c \kappa^{2/3}}{b^{4/3}}$.
- The CLIC drive beam needs a complex system of bunch compression/de-compression to mitigate the effects of CSR.
- The phase stability of the CLIC drive beam is a critical parameter.
- Limited energy spread → Difficult compression/de-compression.
- Shielding could loosen the requirement for compression - accurate simulation is critically needed.
- LHeC has got interest in plate separations small compared to the bunch length.



Schematic representation of the CLIC drive beam complex.

- The implementation of regular CSR in PLACET is based on Saldin et al NIM A 398 (1997) 392.
- Choose an approach based on image charges to match the existing CSR implementation.

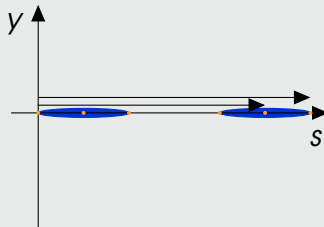
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Normal CSR

- The beam interacts with itself through an electromagnetic field
- Very low energy photons *sim* the minimum wavelength is approximately the bunch length.
- The wake propagates **ahead** of the emitting particle.
- The beam is assumed to have no transverse extent (1 dimension).
- One dimensional model.

CSR shielding

- The beam travels between parallel plates separated by a distance H .
- Like being between two perfectly reflecting mirrors.
- The propagating photons must travel longer to interact.
- \rightarrow The photons can interact with particles in the back of the bunch.
- 1 dimensional model.
- One dimensional condition: $\sigma_x \ll \rho^{1/3} \sigma_z^{2/3}$ must be fulfilled for accurate results.



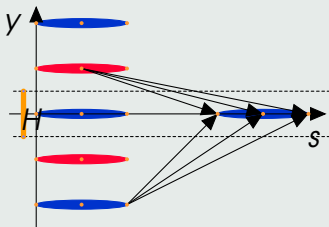
NOTE Different conventions of the orientation of the s -axis. Internally, both normal CSR and the implementation of shielding use

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$$\begin{aligned}
 \left. \frac{d\mathcal{E}_{\text{CSR}}}{ds}(s) \right|_{B_1} = & N r_c m c^2 \left\{ \int_{\alpha_a}^{\alpha_b} d\alpha \left(\frac{\beta^2 \cos(\alpha) - 1}{2|\sin(\alpha/2)|} + \frac{1}{\gamma^2} \frac{\text{sgn}(\alpha) - \beta \cos(\alpha/2)}{\alpha - 2\beta|\sin(\alpha/2)|} \right) \lambda'(s_\alpha) - \frac{\kappa_1 \lambda(s_\alpha)}{2|\sin(\alpha/2)|} \right|_{\alpha_a}^{\alpha_b} \\
 & + \int_{\Delta_a}^{\infty} d\Delta \frac{1}{\gamma^2} \frac{\lambda'(z - \Delta)}{\Delta} + \int_{\Delta_b}^{\infty} d\Delta \frac{1}{\gamma^2} \frac{\lambda'(z + \Delta)}{\Delta} \\
 & + \sum_{n=1}^{\infty} 2(-1)^n \left[\frac{-\kappa_1 \lambda(s_{\alpha,n})}{r_{\alpha,n}} \right|_{\alpha_a}^{\alpha_b} + \int_{\alpha_a}^{\alpha_b} d\alpha \frac{\beta^2 \cos(\alpha) - 1}{r_{\alpha,n}} \lambda'(s_{\alpha,n}) \right] \Big\} \quad (\text{A1})
 \end{aligned}$$

with the definitions

$$\begin{aligned}
 \alpha_a &\equiv \kappa_1(s - B_1), & \alpha_b &\equiv \kappa_1 s, & \Delta_a &\equiv s - 2\beta \frac{1}{\kappa_1} \sin\left(\frac{\kappa_1 s}{2}\right), & \Delta_b &\equiv B_1 - s + 2\beta \frac{1}{\kappa_1} \sin\left(\frac{\kappa_1(B_1 - s)}{2}\right), \\
 r_{\alpha,n} &\equiv \sqrt{2 - 2\cos\alpha + (n\kappa_1 H)^2}, & s_\alpha &\equiv s - s_0 - \frac{1}{\kappa_1}(\alpha - \beta\sqrt{2 - 2\cos\alpha}), & s_{\alpha,n} &\equiv s - s_0 - \frac{1}{\kappa_1}(\alpha - \beta r_{\alpha,n}).
 \end{aligned} \quad (\text{A2})$$

NOT included:

- Transverse effects.
- Reflection of photons on beam pipe.
- 3D extent of bunches.

- C. Mayes and G. Hoffstaetter, Exact 1D model for coherent synchrotron radiation with shielding and bunch compression, PRST-AB **12**, 024401 (2009)
- Beginning principle is Jefimenko form of Maxwell's equation (the usual approach is Lienard-Wiechert fields of relativistic charges)

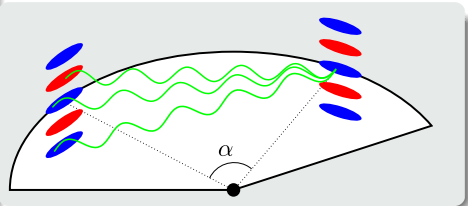
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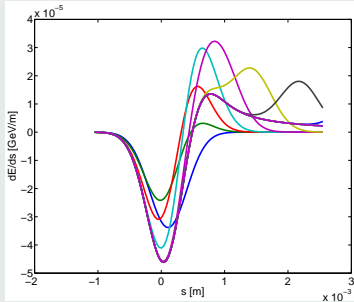
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- Terms 1 and 3 reduce to the CSR already implemented in PLACET when α is small.
- Terms 2,4 and 5 neglected due to $1/\gamma^2$ scaling.
- The (sum of) terms 6 and 7 are CSR shielding. These terms are newly implemented.
- Ultrarelativistic: $\beta = 1$ used.
- Notice the similarity between CSR and CSR shielding.
- Magnitude of wake is energy independant when ultrarelativistic.
- Note that when $H < l_b$, the transient term at α_a is inside the particle distribution.
- Energy conservation is not manifest.

- Choose to perform the integration using the trapezoidal rule.
- Integration is over the retarded angle α \rightarrow allows for accurate inclusion of bunch compression when the bunch distribution is evaluated at the correct *alpha*.
- Inclusion of bunch shape memory (bunch compression) leads to possible decreased numerical stability.
- Numerical stability was re-gained by binning particles differently longitudinally (± 3 sigma from mean value).
- Numerical instabilities are particularly prominent at small plate separations.



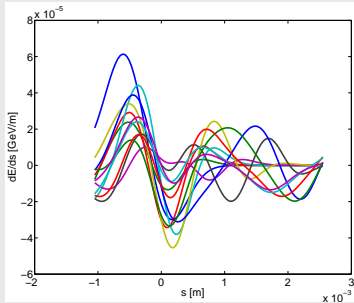


CSR

- The wake varies along the length of the bunch.
- The wake builds up as the magnet is traversed.
- As expected the wake propagates forward and reaches steady state.

CSR shielding

- When image charges are introduced, the wake becomes much more complex.
- As expected the effect vanishes for large plate separations.
- With zero plate distance and 1 image charge, 2 times the normal CSR wake with opposite sign.
- The wake will not reach steady state. New image charges (further away) will always be able to interact (less strongly).



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Savitzky-Golay filtering



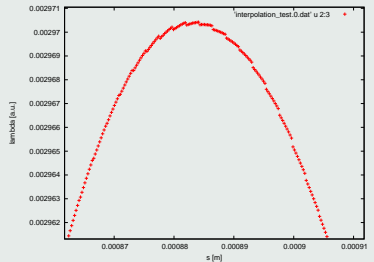
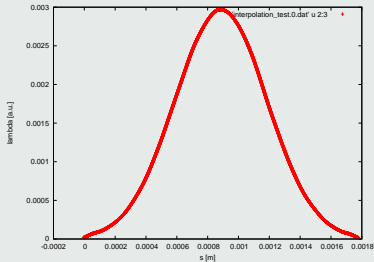
Source: Wikipedia

- Placet already uses Savitzky-Golay filtering to evaluate the charge distribution and its derivative.
- The method does polynomial least-squares fits to a point and a few of its surrounding points - And evaluates the polynomial in the point of interest.
- Normal CSR only needs to evaluate the distribution at bin centers - we would like to evaluate it anywhere.
- Method allows the evaluation of bunch density and its derivative in positions between bins.

Savitzky-Golay interpolation



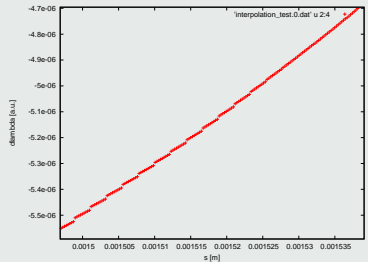
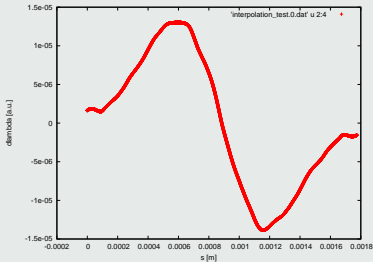
- Since an n 'th order polynomial is available at each point, one can do interpolation to this order.
- Small residual numerical noise from the interpolation.



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New parameters of SBends in PLACET



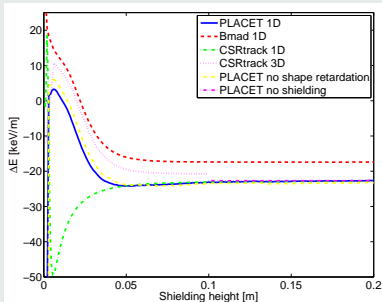
- A new set of parameters are needed to simulate shielding.
 - Some parameters are re-used, but take on an additional role in the shielding case.
 - PLACET input files are scripts written in Tcl, but names of the variables are available in e.g. the Octave interface as well.
-
- Used only by csr shielding.
 - `csr_shielding`: Switch on csr shielding.
 - `csr_shielding_n_images`: Minimum number of images charges (on one side of the plates) used by shielding. This is at the magnet entrance modified at the entrance to the magnet to
$$\text{csr_shielding_n_images} > \frac{R}{H} \sqrt{(\theta + l_b)^2 - 4 \sin^2(\theta/2)}$$
with a warning.
 - `csr_shielding_height` Shielding height (m)
 - Used by both csr without and with shielding.
 - `csr` Switch on csr.
 - `csr_charge` Total bunch charge.
 - `csr_nbins` Number of bins used to evaluate longitudinal distribution and is derivative
 - `csr_filterorder` Order of polynomial used for Savitzky-Golay filtering.
 - `csr_nhalfilter` Number of bins used on either side of a bin for Savitzky-Golay filtering.
 - `csr_nsectors` Number of sectors the magnet is split into ($\propto 1/\Delta s$)

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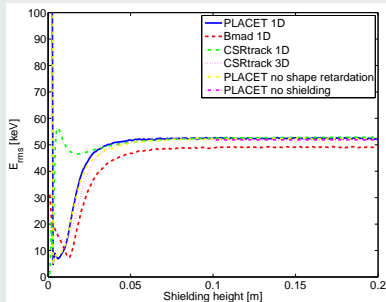
Comparison with other codes



Final mean energy



Final RMS energy

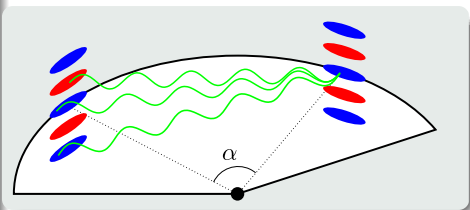


Benchmarking case

- $E_0 = 5\text{GeV}$.
- $L_{mag} = 3\text{m}$
- $\rho = 10\text{m}$
- Bunch length 0.3mm
- 1nC bunch charge
- $4 \cdot 10^5$ macro particles
- 800 CSR bins
- 0.1m step length
- >64 image charges.

Parameter set chosen to match that of Phys. Rev. ST Accel. Beams 12, 040703 (2009)

- None of the codes obey energy conservation at small plate separations.
- Even so, there is a similar **change** in mean+RMS energies when varying the plate distance.
- **Not a full input parameter space optimization has been done for other codes than PLACET.** There could potentially be regions of input parameter space where other codes behave differently.

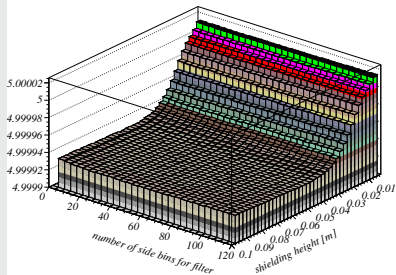


Examples and benchmarking

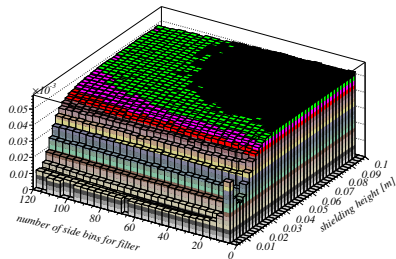


- Select a point where shielding is very strong: $H=6\text{mm}$.
- Choose a set of parameters that seems reasonable: $\text{csr_nbins}=1000$, $\text{csr_nsectors}=100$, $\text{csr_filterorder}=3$, $\text{csr_nhalffilter}=30$.
- Vary parameters around this parameter set.
- Simulations are stable around this point.

Mean energy



RMS energy

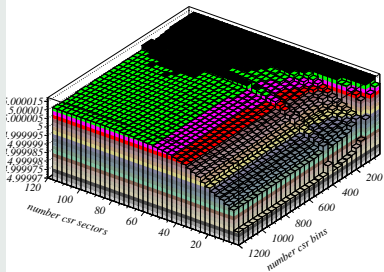


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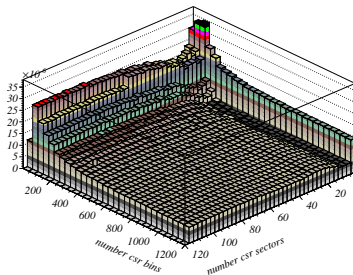


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RMS energy

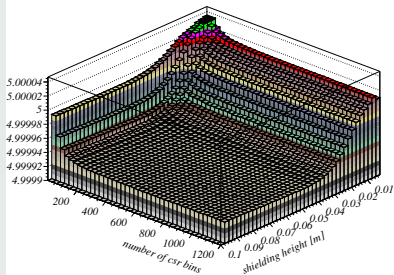


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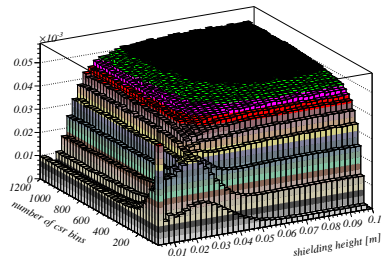


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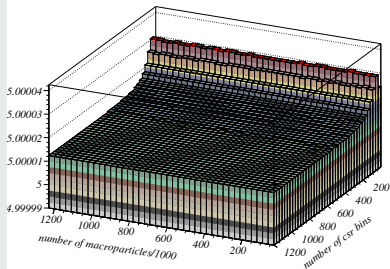


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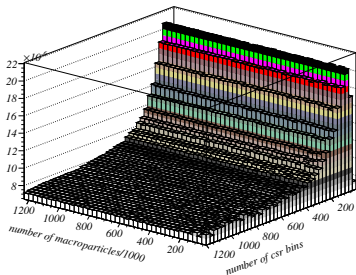


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Mean energy



RMS energy



Phase space of final beam.



Bmad

PLACET

- Without shielding, there is some discrepancy between Bmad and PLACET.
- PLACET with no shielding shows perfect agreement with ELEGANT (E. Adli).
- When decreasing the parallel plate distance, the shielding wake can start to interact with the tail of the bunch.
- Large difference between Bmad and new PLACET implementation for small plate separations.

- CSR shielding is inherently difficult to simulate accurately. Particularly for small parallel plate separations.
- An 1D simulation of CSR shielding with bunch compression has been implemented in PLACET.
- User can input virtually any bunch distribution.
- There are differences between results obtained with different codes.
- Experimental input from e.g. V.Yakimenko et. al Proceedings of 2011 PAC, WEP107, might prove helpful.
- In-house experiments at CTF3 H. H. Braun et. al SLAC-PUB-9353 do not show detailed results on the energy distribution.
- Micro-bunching effects could in principle be simulated, but this needs testing.
- <http://savannah.cern.ch/projects/placet>